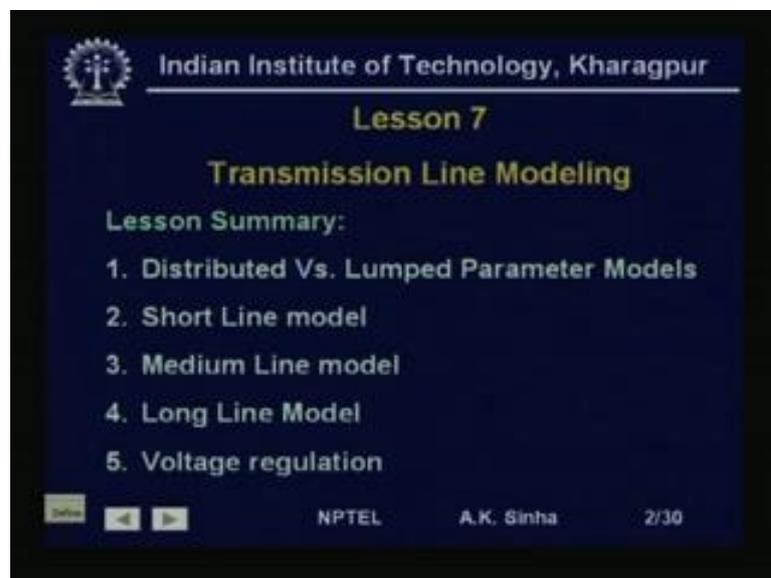


Power System Analysis
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Lecture - 7
Transmission Line Modeling

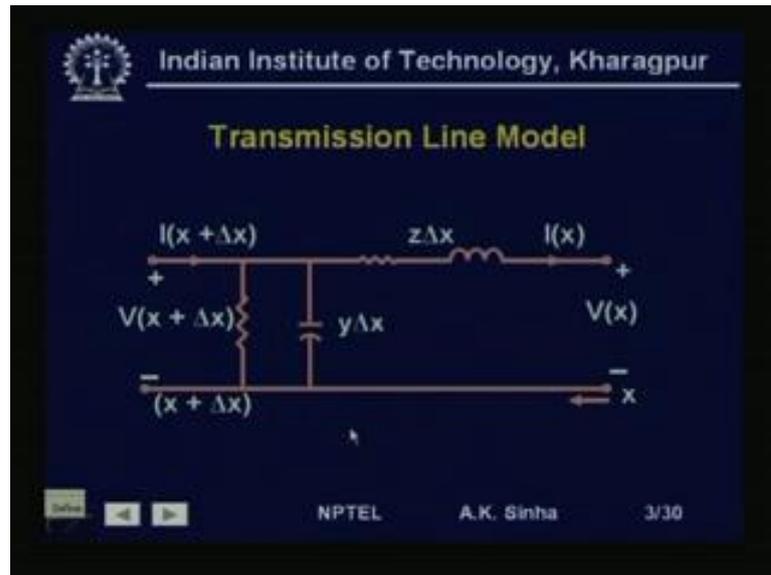
Welcome to lesson 7 in Power System Analysis. In this lesson, we will talk about Transmission Line Modeling.

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In this first we will talk about distributed versus lumped parameter models. Then, we will talk about short line model, medium line model and long line model. We will also discuss voltage regulation for transmission lines. Well, as we have seen in our previous lessons, when we discussed the calculation for transmission line parameters. We found that transmission line have in general three parameters, the line resistance, the line inductance and a capacitance of the line to neutral or ground. Now, as we had seen in the previous lessons, these parameters are distributed all along the line. That is if we see a line it will be something like this, where I have taken a small section of a line.

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For this line of a length Δx , we have the transmission line impedance. Or the series impedance of the transmission line, which is given by z into Δx . It consists of the resistance of the line and the inductance of the line. These will be per unit length or per meter multiplied by Δx will give us the value of series impedance, $z \Delta x$ for the Δx length of the line.

Similarly, we will have a capacitance to ground for the line which will be y ; and for Δx of the line this shunt admittance is going to be y into Δx . There is capacitance will provide us an admittance $j \omega c$, which will be equal to y . And for that will be per unit length of the line multiplying it by Δx the length of the line we get $y \Delta x$. Here I had also shown a conductance part for the line normally for 60 Hertz line or 50 Hertz line the power frequency line.

This conductance is negligible and most of the time we neglect this conductance. So, we have resistance and inductance per unit length of the line. This will form the series impedance per unit length of the line. And we have the capacitance per unit length of the line, which forms the shunt admittance of the line. Now, when we talk about modeling of transmission lines, what we have to see is. If we use distributed parameter model of the line, then the modeling of the transmission line becomes very complex. And it will result into differential equations.

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Lumped Parameter Models

For Sinusoidal waves on O/H lines:

$$\lambda = \frac{C}{f}; C = 3 \times 10^8 \text{ m/s}$$

for $f = 50 \text{ Hz}$; $\lambda = 6000 \text{ km}$

If the line length is $< 250 \text{ km}$ long, then one can neglect the distributed effect of the line parameters and it is sufficient to consider lumped parameter model for transmission lines

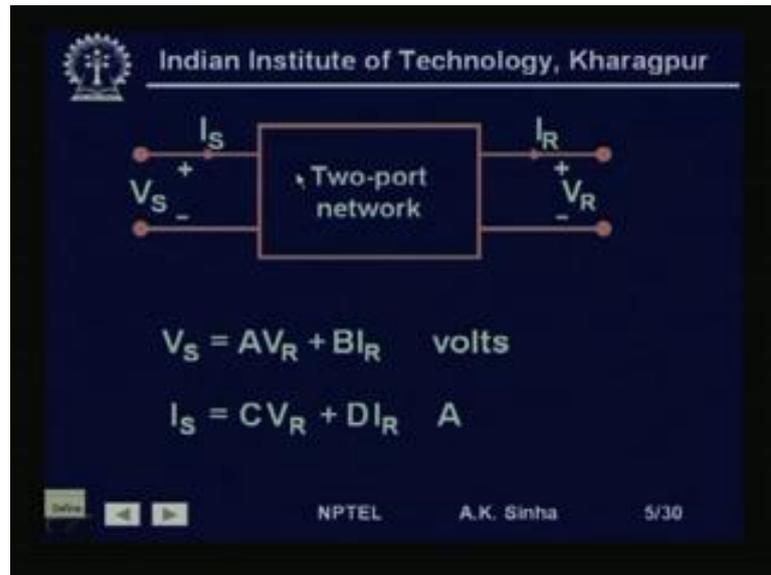
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What we will see is that for sinusoidal voltage and current waves on overhead line. The wavelength of the line is given by this lambda; wavelength of the line is given by C by f , where, C is the velocity of light or the velocity of propagation of the waves, current or voltage waves on the transmission line. For overhead transmission line, this will be very nearly equal to the velocity of light.

Therefore, wavelength lambda is equal to C by f , where C is 3 into 10 to the power 8 meters per second. If we substitute this value of C and for frequency of 50 hertz, we get lambda is equal to 6000 kilometers. That is the wavelength of a transmission line at 50 Hertz is of the order of 6000 kilometers. Since most of the transmission lines that we have, are only a few 100 kilometer length.

Therefore, in most of the cases, we can neglect the distributed parameter model of the line. That is the distributed effect can be neglected and it is sufficient to model the line as lumped parameter model. That is if line length is less than about 250 kilometers, then one can neglect the distributed effect of the line parameters. And it is sufficient to consider lumped parameter model for transmission lines. Now, once we have used a lumped parameter model for the transmission line. Then, we can consider this transmission line as a simple two-port network.

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Where we have an input port, where we have the voltage V_S and I_S . That is the sending end port, V_S is the sending end voltage, I_S is the sending end current. And V_R is the receiving end voltage and I_R is receiving end current. So, transmission line can be visualized as a two-port network. For this two-port network, we can write down the relationship between the sending and voltage and current. In terms of receiving and voltage and current as V_S is equal to A times V_R plus B times I_R volts.

And I_S the sending and current is equals to C times V_R , that is receiving an voltage. Plus D times I_R the receiving in a current. Of course the unit will be amplic. Now, this is a general relationship for two-port network. It is valid for any two passive bilateral linear two-port network. And one of the properties of this from network theory, we get one of the properties for this network as $AD - BC$ is equal to 1. That is if we are putting this relationship in a matrix form, we get V_S, I_S is equal to $ABCD$ into V_R, I_R . This is the general transmission line ABCD parameter model for any transmission line with the relationship that $AD - BC$ is equal to 1.

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Short Line Model

- Line Length < 80 km
- Generally MV / LV Lines
- Capacitance can be neglected

$Z = z l = (R + j \omega L) l$

I_S I_R

V_S V_R

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Now, transmission lines which are less than 80 kilometers long or about 50 miles long, we neglect the effect of capacitance. Because, the capacitive effect or the charging current is not very large and its effect can be neglected. Also these lines are generally medium voltage and low voltage lines, because these are very short lines. And therefore, also the charging current, because the voltage being small will also be much smaller.

So, generally for line lengths less than 80 kilometers, which are generally medium voltage or low voltage lines, sometimes even high voltage lines can be there. The capacitance for these lines can be neglected, in that case the lumped parameter model for the transmission line, will consist of only the resistance and inductance. That is the series impedance of the line, the series resistance of the line and the series inductance of the line.

Here, the series impedance of the line Z is equal to z into l , small z into l , where small z is the per unit length series impedance. That is R plus $j \omega l$ per unit length, that is resistance per unit length and inductance per unit length, as we have seen we can calculate it for any line. So, Z is equal to z into l , this is equal to R plus $j \omega l$ into l , where R is the resistance per unit length of the line and $j \omega l$ is the reactance per unit length of the line.

So, from this model we have V_S the sending end voltage, I_S the sending end current only series impedance, that is here R and this $j \omega l$ or X . And I_R is the receiving

end current and V_R is the receiving end voltage. From this circuit it is very clear I_S and I_R will be equal. Therefore, we can write down the circuit equation for this transmission line model as V_S is equal to V_R plus Z into I_R .

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$$V_S = V_R + Z I_R$$

$$I_S = I_R$$

$$\begin{bmatrix} V_S \\ I_S \end{bmatrix} = \begin{bmatrix} 1 & Z \\ 0 & 1 \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

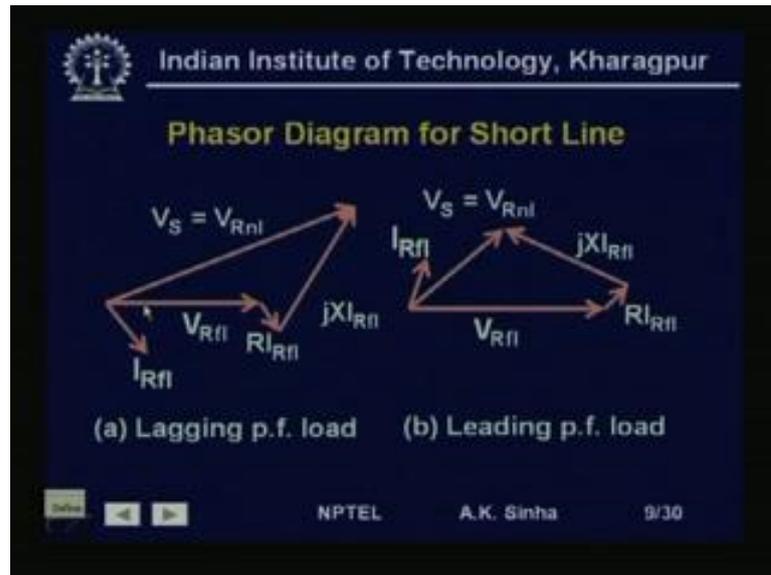
A = D = 1 per unit
 B = Z Ω
 C = 0 S

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V_S is going to be equal to this voltage plus the drop in this transmission line impedance. Series impedance of the transmission line, this will be equal to Z into I_R . So, we have V_S is equal to V_R plus Z into I_R . And as we have all ready seen I_S will be equal to I_R , therefore writing it in terms of matrix equation for ABCD parameters. We have V_S I_S is equal to V_R plus Z into I_R and I_S is equal to $0 V_R$ into plus 1 into I_R . That is ABCD parameters from here, if you see are A is equal to D is equal to 1 per unit.

And B is equal to Z ohms and C is equal to 0 Siemens, that is the unit of A and D is basically dimensionless, it will be in terms of per unit. Whereas, the unit of B is the unit of impedance or ohms and unit of C is the unit of admittance or Siemens. So, from this relationship for a short line, we have A and D is equal to 1 B is equal to Z and C is equal to 0.

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The Phasor diagram for short line can be developed, we can start with the voltage at the receiving end at full load. So, V_{Rfl} as the reference voltage and we have a current I_{Rfl} at full load, which is in this direction, that is lagging the voltage by some angle θ . Now, with this we can find out using the circuit model, voltage at the sending end, which will be equal to V_{Rfl} plus $R I_{Rfl}$.

That is resistance, the drop in the resistance due to the current, receiving end current at full load. Plus the drop in the reactance, that is $jX I_{Rfl}$, so the drop in the reactance. So, when these drops are added, we get the voltage V_S . That is what we see is for lagging power factor load, where I_{Rfl} is lagging V_{Rfl} by an angle θ . We get V_S which will generally be larger than V_{Rfl} , that is there is going to be a positive drop of voltage for the system.

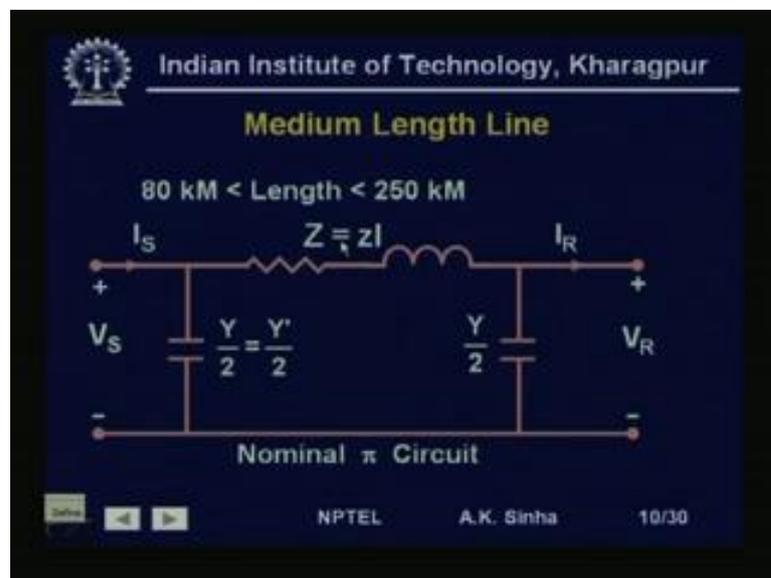
Similarly, if the current is of leading power factor, that is again choosing V_{Rfl} as a reference voltage. And since, the current I_{Rfl} is a leading power factor load. So, it leads this voltage by some angle θ . Then, again by adding the voltage drops, we will get $R I_{Rfl}$, which is the drop in the resistance plus $jX I_{Rfl}$, which is the drop in the reactance. Or the inductive reactance this will give us the value V_S , and this V_S as we see will be many times smaller than V_{Rfl} .

That is in case of leading power factor load, the sending end voltage magnitude can be lowered than the receiving end voltage magnitude. Of course, in case of short lines the V

V_S is also equal to V_R at no load. Because, if there is no current flowing in this, if there is no current flowing in this ((Refer Time: 15:29)), then the voltage V_R and V_S will be equal. So, V_R at no load is equal to V_S , that is what we have shown here. That is V_S is equal to V_R at no load V_R n l, this we required for calculating regulation as we will see later.

Now, for line lengths or lines having a length between 80 to 250 kilometers, we call these lines as medium Lang lines. These lines are generally high voltage or extra high voltage lines. And for these lines, we cannot neglect the shunt capacitance or the capacitive admittance of the line. The shunt admittance needs to be taken care of.

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What we do here is, again we are use a lumped parameter model where the complete series impedance that it is total resistance of the line, is lumped as one resistance. And the total inductance of the line is again lumped together. So, we get this R plus j ωL as the series impedance of the line. So, Z the total series impedance is equal to z into l , where small z is the series impedance of the line per unit length.

Now, the total shunt admittance of the line will be Y , which will be equal to small y the shunt admittance of the line per unit length into l . Now, what we do is normally, we divide this total shunt admittance into two parts and put them at the two ends. So, we have shunt admittance Y by 2 , put at the sending end and another Y by 2 put at the receiving end.

So, this model it looks like a pi model pi and therefore, we call the model as a nominal pi circuit for medium length line. Now, again here we have this as a two-port network with V_S, I_S the sending end voltage end current and V_R, I_R the receiving end voltage end current.

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$$V_S = V_R + Z \left(I_R + \frac{V_R Y}{2} \right)$$

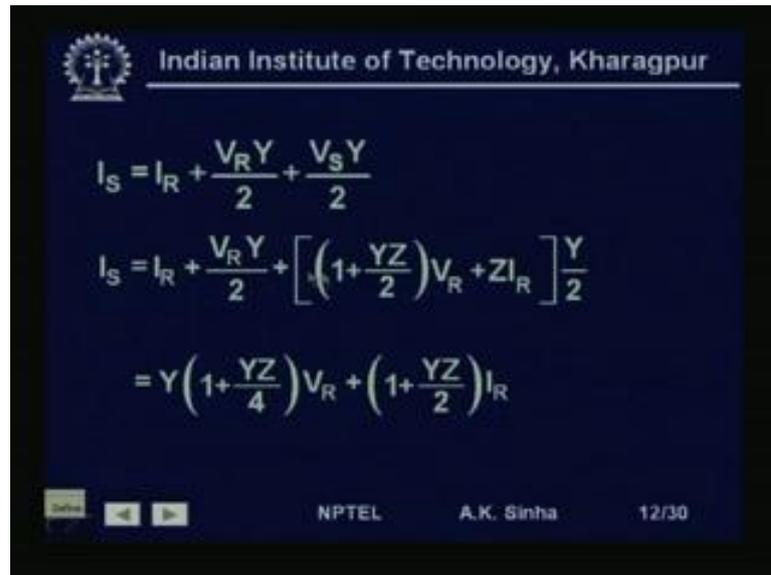
$$= \left(1 + \frac{YZ}{2} \right) V_R + Z I_R$$

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And we can use the ABCD parameter for this as V_S , from here we can see is equal to V_R plus Z into I_R plus V_R into Y by 2. Now, if you look at this circuit, what we are finding is, V_S this voltage will be equal to the drop in this. Now, what will be the drop in this, this drop will be equal to how much current is flowing in this part I_R plus V_R into Y by 2.

So, I_R plus V_R into Y by 2 is the current this multiplied by Z will give us the drop in this. So, that is what we are writing V_R is the receiving end voltage plus the drop, drop is Z into I_R plus V_R Y by 2. So, this is the current which is flowing in the charging admittance of the line placed at the receiving end. So, V_S is equal to V_R plus Z into I_R plus V_R Y by 2. If we rearrange this by combining all the terms of V_R together, we get this as 1 plus $Y Z$ by 2 into V_R plus Z into I_R . Similarly, we can write down the relationship for the sending end current. Now, sending end current sending end current I_S is going to be equal to this I_R plus the current here. That is V_S into Y by 2 plus the current here which is V_S into Y by 2.

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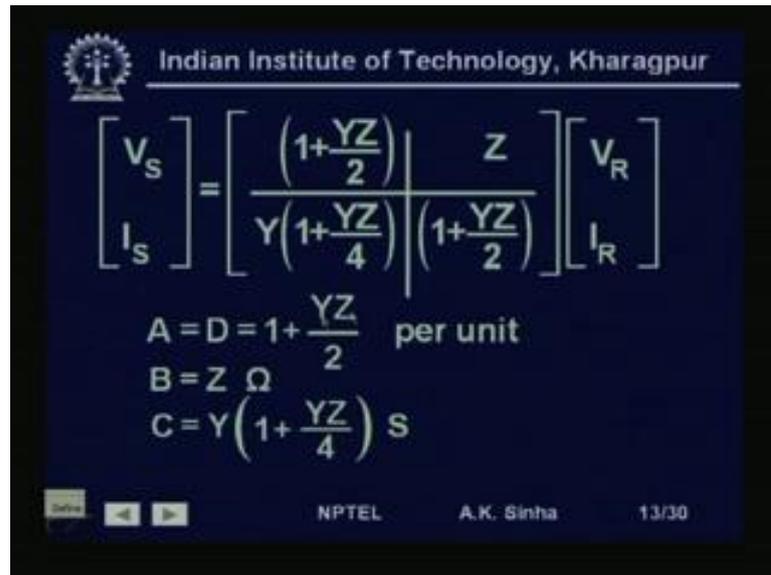
$$I_S = I_R + \frac{V_R Y}{2} + \frac{V_S Y}{2}$$
$$I_S = I_R + \frac{V_R Y}{2} + \left[\left(1 + \frac{YZ}{2} \right) V_R + Z I_R \right] \frac{Y}{2}$$
$$= Y \left(1 + \frac{YZ}{4} \right) V_R + \left(1 + \frac{YZ}{2} \right) I_R$$

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So, we are writing I_S is equal to I_R plus V_R into Y by 2 plus V_S into Y by 2 . Now, again rearranging this by combining terms, we write this as I_R plus V_R into Y by 2 plus V_S . We are writing the relationship for V_S from the previous equation here, V_S is equal to 1 plus YZ by 2 into V_R plus $Z I_R$. So, substituting this value for V_S , we are putting it here, so V_S into Y by 2 we write here.

So, this is 1 plus YZ by 2 V_R plus $Z I_R$ for V_S into Y by 2 . Now, again rearranging it, we will get this as Y into 1 plus YZ by 4 into V_R plus 1 plus YZ by 2 into I_R . This is the relationship for the current I_S , in terms of V_R and I_R . So, now using these relationships, we can write down this.

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$$\begin{bmatrix} V_S \\ I_S \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{YZ}{2}\right) & Z \\ Y\left(1 + \frac{YZ}{4}\right) & \left(1 + \frac{YZ}{2}\right) \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

$A = D = 1 + \frac{YZ}{2}$ per unit
 $B = Z \Omega$
 $C = Y\left(1 + \frac{YZ}{4}\right) S$

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In terms of ABCD parameters as V_S is equal to $1 + \frac{YZ}{2} V_R$ plus $Z I_R$, I_S is equal to $Y\left(1 + \frac{YZ}{4}\right) V_R$ plus $\left(1 + \frac{YZ}{2}\right) I_R$, which shows that A is equal to D , that is A is equal to D is equal to $1 + \frac{YZ}{2}$ this will be in π unit. B will be equal to Z ohms and C will be equal to $Y\left(1 + \frac{YZ}{4}\right)$ Siemens. So, for a medium length line, when we are using a nominal π equivalent for this medium length line, we have the ABCD parameters.

Given as A is equal to D is equal to $1 + \frac{YZ}{2}$, B is equal to Z ohms and C is equal to $Y\left(1 + \frac{YZ}{4}\right)$ Siemens. Now, we will talk about voltage regulation. Now, voltage regulation, what we mean by voltage regulation is basically in when there is no load on the line, the voltage on the line will be somewhat higher. When we load the line, generally the loads will be of a lagging power factor loads.

So, when we load these lines, then the voltage at the receiving end drops. And voltage regulation is basically telling us how much this voltage drop is, in terms of the rated voltage or the sending end voltage. Therefore, we define this voltage regulation as a percentage voltage regulation.

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Voltage Regulation

$$\text{percent VR} = \frac{|V_{Rnl}| - |V_{Rfl}|}{|V_{Rfl}|} \times 100$$

For Short Line;
 $V_{Rnl} = V_s$

For Medium Length Line;
 $V_{Rnl} = AV_s$

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So, percent voltage regulation is defined as V_R no load minus V_R full load divided by V_R full load into 100 percent. So, this is what we define, basically what we are trying to say here is that suppose we have a line which is loaded to it is full load value. And if the load is suddenly thrown off what will be the voltage at the receiving end. How much rise of voltage is there at the receiving end.

So, this V_{Rnl} as showing what will be the voltage of the receiving end? When the load is thrown off and V_{Rfl} is the voltage when the full load is there on the transmission line. So, this percentage voltage regulation gives us a very good parameter. Because, when we are designing a transmission line, we cannot allow more than about 8 to 10 percent voltage drop.

So, voltage regulation will tell us what is going to be the voltage drop, when the transmission line is fully loaded, because, if the voltage drop is more, then we need some compensation. We will also we had seen that when the load is capacitive in nature or of leading power factor, then the receiving end voltage can go higher than the sending end voltage. And this also needs to be checked for the transmission line.

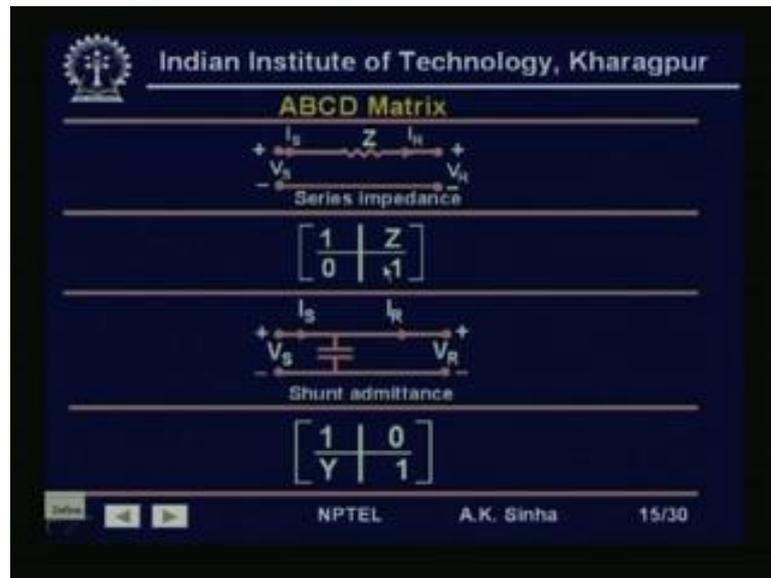
Because, if the receiving end voltage becomes very high, then it may endanger the transmission line insulators, because the design voltage will be based on the nominal value or the rated value for the transmission line. Therefore, if you see here for the short line, when we want to calculate the voltage regulation we need to calculate the voltage at

the receiving end at no load. And as we have seen earlier the voltage at the receiving end at no load, for the short the line is nothing but, equal to V_S .

That is when I_R is equal to 0, that is I_R and I_S both are equal to 0, because I_R there is only I_S and I_R are equal. So, when the load current is 0, that is the system is at no load, then this drop is also equal to 0, so V_R is equal to V_S . For a medium length line, that is the pi model that we have used, we have seen that V_S is equal to $A V_R$ plus Z into I_R . Now, if we see this then when the I_R is 0, that is at no load we have V_S is equal to $A V_R$ plus Z into 0. So, V_S will be equal to $A V_R$ only.

And therefore, we have the no load voltage at the receiving end, that is V_R at no load will be equal to A into V_S . And therefore, we can calculate V_R at no load from this relationship. Now, we will take up the ABCD matrix for various kinds of line configurations. That is various kinds of network models or parameters that we use.

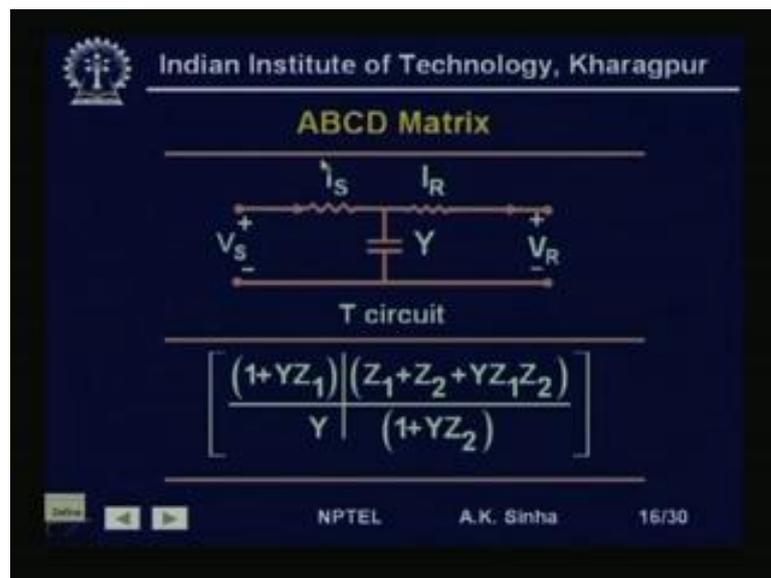
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So, if we have a simple series impedance which is the case for a short line. Then we have the ABCD parameters as A is equal to 1, B is equal to Z , C is equal to 0 and D is equal to 1. If we have a simple capacitance, this is what we many times use when we have to compensate for the reactive power drawn by the load. So, many times when the load which is highly inductive, draws a large amount of reactive power. We need to compensate this reactive power by putting capacitance at the receiving end.

The shunt capacitance at the receiving ends, will then provide some reactive power. And this will compensate for the large amount of reactive power drawn by the load or consumed by the load. So, capacitance or the capacitor will generate reactive power at the receiving end itself. And compensate for the large amount of reactive power drawn by the load. For simple shunt capacitance the ABCD parameters will be A is equal to 1, B is equal to 0, C is Y and D is equal to Y.

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Similarly, instead of using a pi equivalent model, if you use a T equivalent model, where what we do is we divide the total series impedance into two equal parts. And lump the total capacitance as one capacitance and put it at the centre of the line. Then in this model we have Z by 2 and Z by 2 here, and total Y is here, in this kind of a model which is a T equivalent circuit for the transmission line. The ABCD parameters will be given by A is equal to $1 + YZ_1$, where this is Z_1 and this is Z_2 .

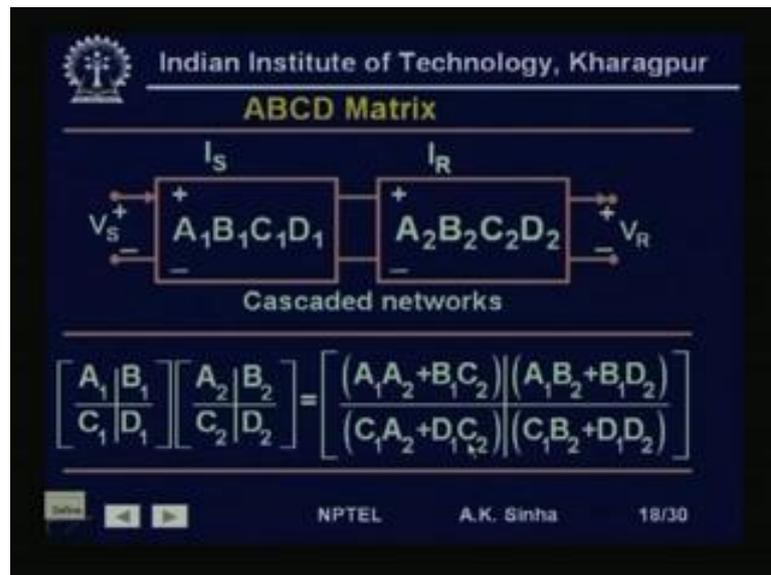
So, $1 + YZ_1$ and B will be equal to $Z_1 + Z_2 + YZ_1Z_2$. So, and C will be equal to Y and D will be equal to $1 + YZ_2$. Again for this if you see $AD - BC$ will be equal to 1, for the pi equivalent circuit, we have all ready seen. Now, here for this pi circuit we have Y_1 and Y_2 which may not be equal, whereas for the nominal pi circuit we had both $Y_1 = Y_2$ where equal.

That is the total shunt capacitance was divided equally and placed half of it was placed at the sending end and half was placed at the receiving end. Now, for this general pi circuit,

if we calculate the ABCD parameters, we will get A is equal to $Y_1 + Y_2 Z$, B will be equal to Z, C will be equal to $Y_1 + Y_2 + Y_1 Y_2 Z$ and D will be equal to $1 + Y_1 Z$. Again here if you see $AD - BC$ will be equal to 1.

Now, as I said many times we use compensation like adding a capacitance at the sending end. And then the line model will consist of two parts, one will be the transmission line model, another will be for the compensating equipment.

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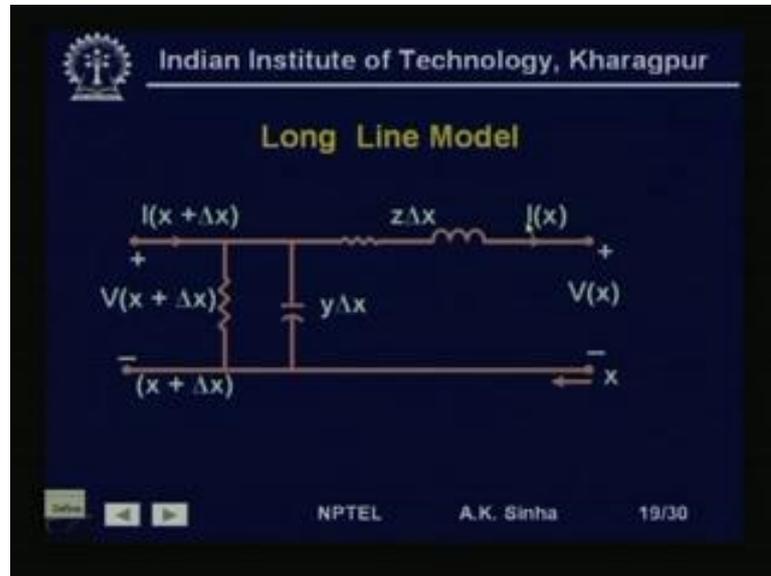


And in that case we have two cascade ABCD systems. And for this complete system also we can find out the ABCD parameters, which will be nothing but, this $A_1 B_1 C_1 D_1$ multiplied by $A_2 B_2 C_2 D_2$ will give the composite ABCD of this. That is A for the composite system or cascaded network will be $A_1 A_2 + B_1 C_2$. B will be equal to $A_1 B_2 + B_1 D_2$, C will be equal to $C_1 A_2 + D_1 C_2$ and D will be equal to $C_1 B_2 + D_1 D_2$.

So, in this way we had seen how we can model the transmission line, for short lengths and medium lengths. But, when transmission line lengths are very long, which is the case in extra high voltage lines. Or some of the lines which are connecting the remote power generating sources, such as hydel plants to the load centered. Then in that case it is no longer possible to use a lumped parameter model.

And we will have to use the distributed parameter model, or the more accurate model otherwise. The errors which will come, because of using the lumped parameter model may not justify the simplification that we have used.

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So, here again we come back to the distributed parameter model. So, for long line as I said, we will use this distributed parameter model. We have taken a section of line which is Δx in length a small length Δx . The series impedance of the line is z into Δx , where z is the series impedance of the line the resistance plus the inductive reactance of the line R plus $j\omega l$ per unit length. So, that is represented here as z into Δx .

And the total shunt admittance is y into Δx , for this Δx part of the line where small y is the shunt admittance of the line per unit length. As I said for the transmission line at power frequency 50 Hertz or 60 Hertz, this part the conductance is negligible and this is generally neglected and only capacitance is there. Now, we have taken the parameter x or the distance parameter from the receiving end side.

So, the receiving end voltage or the voltage at a point at a distance x from the receiving end we write this voltage as V at distance $V x$. And the current at this point as I at a distance x from the receiving end that is $I x$. Similarly, the voltage of the line Δx meter in this positive direction, that is from receiving end towards the sending end will be given by V into x plus Δx . That is voltage at x plus Δx from the receiving end; current will be given by I at x plus Δx , that is at current at x plus Δx distance

from the receiving end. Now, with these values now we will write the Kirchhoff's voltage law.

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$$z = R + j\omega L \quad \Omega/m$$

$$y = G + j\omega C \quad S/m$$

$$V(x + \Delta x) = V(x) + (z\Delta x)I(x) \text{ volts}$$

$$\frac{V(x + \Delta x) - V(x)}{\Delta x} = zI(x)$$

$$\frac{dV(x)}{dx} = zI(x)$$

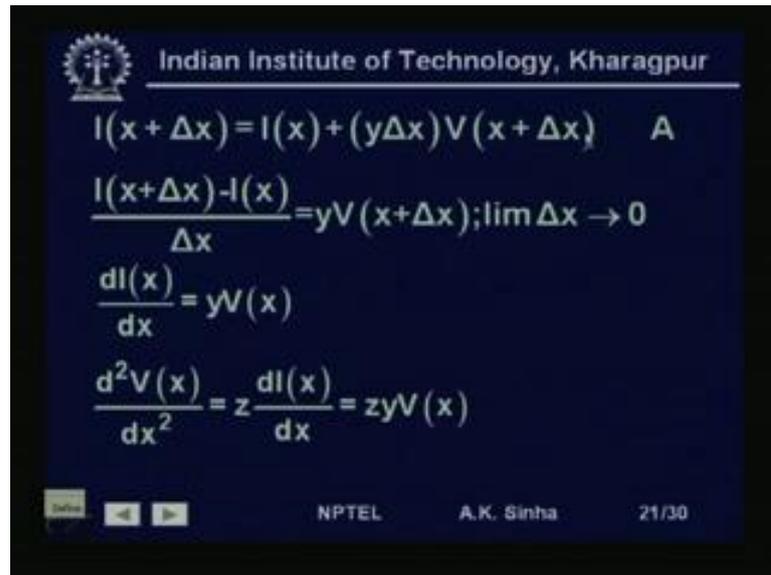
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Here, we have V plus V at x plus Δx will be equal to V at x . That is the voltage at distance x plus the voltage drop, that is the voltage drop here will be I at x into z Δx . So, the voltage at this end, that is V at x plus Δx is going to be equal to V at x plus the voltage drop, which is I at x into z Δx . That is what we have written here, V at x plus Δx is equal to V at x plus z into Δx into I at x .

Now, here as I said earlier z is the series impedance per unit length, which will be R plus j ω L ohms per meter. And y is G plus j ω C Siemens per meter, where G for power frequency or transmission lines we have is almost 0 or negligible. Now, from this relationship of voltage at distance x plus Δx , in terms of voltage at distance x , if you rearrange, this we will get as V at x plus Δx minus V at x divided by Δx .

This Δx term we have taken on this side and this V at x , we have taken on this side, so this becomes negative. And this Δx comes in the denominator. So, V at x plus Δx minus V at x divided by Δx is equal to z into I at x . Now, if we take a limit with Δx tending to 0, then we can write this as $dV(x)$ by dx is equal to $zI(x)$. That is derivative of voltage at x will be equal to z into I at x .

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$$I(x + \Delta x) = I(x) + (y\Delta x)V(x + \Delta x) \quad A$$
$$\frac{I(x + \Delta x) - I(x)}{\Delta x} = yV(x + \Delta x); \lim_{\Delta x \rightarrow 0}$$
$$\frac{dI(x)}{dx} = yV(x)$$
$$\frac{d^2V(x)}{dx^2} = z \frac{dI(x)}{dx} = zyV(x)$$

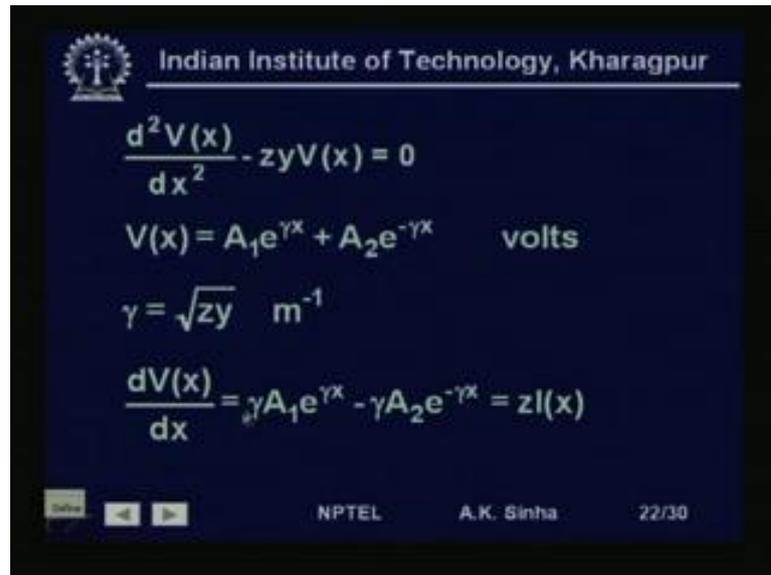
NPTEL A.K. Sinha 21/30

Similarly, writing the Kirchhoff's current law. We will have I at x plus Δx will be equal to I at x plus y into Δx into V at x plus Δx . That is if we look at this ((Refer Time: 38:27)) I at x plus Δx is going to be equal to I at x . Plus the current which is flowing in this admittance which will be equal to y into Δx multiplied by voltage at this point which is V at x plus Δx .

That is why we have got I at x plus Δx is equal to I at x , the current at x plus y into Δx the admittance multiplied by V at x plus Δx . Again rearranging by taking this I at x on this side and dividing by Δx , we have I at x plus Δx minus I at x divided by Δx is equal to y into V at x plus Δx . Now, taking limit Δx tending to 0, we can write this as equal to $dI(x)$ by diagnosis, that is first derivative of current at x with respect to x is equal to y into V at x .

Now, what we have is if we see our previous equation for V . If we differentiate it again with respect to x , then what we will get is $d^2V(x)$ by dx^2 is equal to $z \frac{dI(x)}{dx}$, $d^2V(x)$ by dx^2 is equal to $z \frac{dI(x)}{dx}$. Now, substituting the value of $\frac{dI(x)}{dx}$ from this point or this equation, we get this is equal to z into y into $V(x)$. We can write this as $d^2V(x)$ by dx^2 minus $zyV(x)$ is equal to 0. That is what we have done is taken this term on this side. So, $d^2V(x)$ by dx^2 minus $zyV(x)$ is equal to 0.

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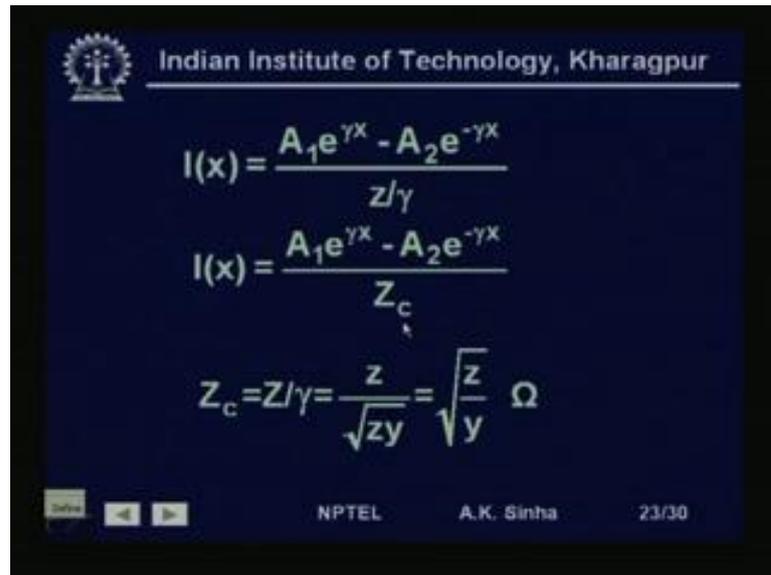
$$\frac{d^2V(x)}{dx^2} - zyV(x) = 0$$
$$V(x) = A_1e^{\gamma x} + A_2e^{-\gamma x} \quad \text{volts}$$
$$\gamma = \sqrt{zy} \quad \text{m}^{-1}$$
$$\frac{dV(x)}{dx} = \gamma A_1e^{\gamma x} - \gamma A_2e^{-\gamma x} = zI(x)$$

NPTEL A.K. Sinha 22/30

Now, this is a second order homogeneous equation. And we know the solution for this equation can be obtained in this form, where $V(x)$ will be given by $A_1 e^{\gamma x} + A_2 e^{-\gamma x}$ volts. So, this will be a general solution for this second order differential equation, where γ is equal to \sqrt{zy} and the unit for this will be per meter.

Now, again differentiating this equation, we can write $dV(x)/dx$ by differentiating this equation with respect to x . We have got $dV(x)/dx$ is equal to $\gamma A_1 e^{\gamma x} - \gamma A_2 e^{-\gamma x}$. And this will be equal to $zI(x)$, because $dV(x)/dx$ is $zI(x)$. Therefore, we can write that $I(x)$ is equal to this term divided by z .

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$$I(x) = \frac{A_1 e^{\gamma x} - A_2 e^{-\gamma x}}{z/\gamma}$$
$$I(x) = \frac{A_1 e^{\gamma x} - A_2 e^{-\gamma x}}{Z_c}$$
$$Z_c = Z/\gamma = \frac{z}{\sqrt{zy}} = \sqrt{\frac{z}{y}} \Omega$$

NPTEL A.K. Sinha 23/30

So, $I(x)$ is equal to $A_1 e^{\gamma x} - A_2 e^{-\gamma x}$ divided by z . And the γ which was here multiplied with A_1 and A_2 that we have taken on in the denominator, so this is z divided by γ . Therefore, we can write $I(x)$ is equal to $A_1 e^{\gamma x} - A_2 e^{-\gamma x}$ divided by Z_c . Now, this term z by γ we are writing this as Z_c , Z_c is called a characteristic impedance of the transmission line.

Now, if you see Z_c as we have put here is z by γ . This is equal to z divided by γ is square root of $z y$ and therefore, this is equal to square root of z by y . That is square root of series impedance by shunt admittance. And the unit for this is ohms and therefore, this Z_c is termed as the characteristic impedance of a transmission line. Now, we need to find out the constants of integration that is this A_1 and A_2 . So, for this what we do is, we use the boundary conditions.

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$$V_R = V(0)$$
$$I_R = I(0)$$
$$V_R = A_1 + A_2$$
$$I_R = \frac{A_1 - A_2}{Z_c}$$

NPTEL A.K. Sinha 24/30

So, here using the boundary, that is at x is equal to 0, V_x is equal to V_R . So, at 0 V at 0 is equal to V_R and I at 0 is equal to I_R , therefore V_R will be equal to A_1 plus A_2 . That is if we substitute in this relationship x is equal to 0, then we get V_0 which is equal to V_R is equal to A_1 plus A_2 as these terms will become 1. So, V_R is equal to A_1 plus A_2 and again substituting it in this ((Refer Time: 44:36)) I_x here x is equal to 0 in these we will get I_x is equal to A_1 minus A_2 by Z_c , A_1 minus I_R is equal to A_1 minus A_2 by Z_c . From these two relationships A_1 plus A_2 is equal to V_R and A_1 minus A_2 by Z_c is equal to I_R , we can calculate A_1 and A_2 .

(Refer Slide Time: 45:06)

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$$A_1 = \frac{V_R + Z_c I_R}{2}$$
$$A_2 = \frac{V_R - Z_c I_R}{2}$$
$$V(x) = \left(\frac{V_R + Z_c I_R}{2} \right) e^{\gamma x} + \left(\frac{V_R - Z_c I_R}{2} \right) e^{-\gamma x}$$

NPTEL A.K. Sinha 25/30

And A 1 comes out to be equal to V_R plus Z_c into I_R by 2 and A 2 will come out to be equal to V_R minus Z_c into I_R by 2. Now, substituting the values of this A 1 and A 2, in the relationship for V_x , we can get V_x is equal to V_R plus Z_c into I_R by 2. That is A 1 into e to the power γx plus V_R minus Z_c into I_R by 2, this is A 2 into e to the power minus γx .

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$$I(x) = \left(\frac{V_R + Z_c I_R}{2Z_c} \right) e^{\gamma x} - \left(\frac{V_R - Z_c I_R}{2Z_c} \right) e^{-\gamma x}$$

$$V(x) = \left(\frac{e^{\gamma x} + e^{-\gamma x}}{2} \right) V_R + Z_c \left(\frac{e^{\gamma x} - e^{-\gamma x}}{2} \right) I_R$$

$$I(x) = \frac{1}{Z_c} \left(\frac{e^{\gamma x} - e^{-\gamma x}}{2} \right) V_R + \left(\frac{e^{\gamma x} + e^{-\gamma x}}{2} \right) I_R$$

NPTEL A.K. Sinha 26/30

And similarly putting it for the value of A 1 A 2 for I_x relationship, we get I_x is equal to V_R plus Z_c into I_R by 2 into Z_c e to the power γx minus V_R minus Z_c I_R divided by 2 into Z_c into e to the power minus γx . Now, if we again rearrange the equations for V_x and I_x , that is in terms of V_R and I_R , then combining all the terms of V_R together and terms of I_R together. Then we can write from this equation take all the terms which are having V_R together and take all the terms which are having I_R together.

Then, we can write V_x is equal to e to the power γx plus e to the power minus γx divided by 2 into V_R plus Z_c e to the power γx minus e to the power minus γx by 2 into I_R . And I_x will be equal to 1 by Z_c into e to the power γx minus e to the power minus γx by 2 V_R plus e to the power γx plus e to the power minus γx by 2 into I_R , this is the relationship that we will get.

Now, here we have got a relationship of V at a distance x, in terms of receiving end voltage and current. Similarly, current at any distance x, in terms of receiving end voltage and current.

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$$V(x) = \cosh(\gamma x)V_R + Z_c \sinh(\gamma x)I_R$$

$$I(x) = \frac{1}{Z_c} \sinh(\gamma x)V_R + \cosh(\gamma x)I_R$$

$$\begin{bmatrix} V(x) \\ I(x) \end{bmatrix} = \begin{bmatrix} A(x) & B(x) \\ C(x) & D(x) \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

NPTEL A.K. Sinha 27/30

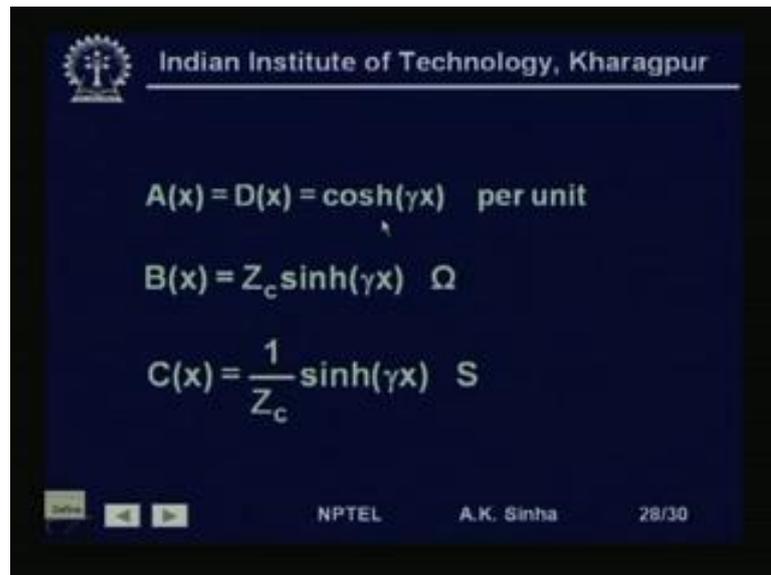
Now, from this we can ((Refer Time: 47:46)) this term $e^{\gamma x}$ plus $e^{-\gamma x}$ by 2 is equal to $\cosh \gamma x$. And $e^{\gamma x}$ minus $e^{-\gamma x}$ by 2 will be equal to $\sinh \gamma x$. Therefore, we have this $\cosh \gamma x$ into V_R plus Z_c into $\sinh \gamma x$ into I_R , that is what we get here.

That is voltage at any point in the transmission line can be found, in terms of the hyperbolic functions and the receiving end voltage and currents. Similarly, the current at any point in the transmission line any distance x from the receiving end, can be found out. Similarly, using the relationship $I(x)$ is equal to $\frac{1}{Z_c} \sinh \gamma x$ into V_R plus $\cosh \gamma x$ into I_R .

Now, if you look at this, this is giving us a relationship of voltage and current, at any point x, in terms of receiving end voltage and current. Therefore, we can write this in terms of ABCD parameters as $V(x) I(x)$ is equal to $A(x) B(x) C(x) D(x) V_R I_R$, where ABCD, A is equal to $\cosh \gamma x$, B is equal to $Z_c \sinh \gamma x$, C is equal to $\frac{1}{Z_c} \sinh \gamma x$. And D is equal to $\cosh \gamma x$, again you will see $AD - BC$ is equal to 1.

Now, since most of the time we are interested only in the terminal conditions. That is conditions at the sending end and the receiving end, not anywhere in between the lines. So, we need to just substitute l instead of x and then we will get the sending end voltage and sending end current.

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$$A(x) = D(x) = \cosh(\gamma x) \text{ per unit}$$

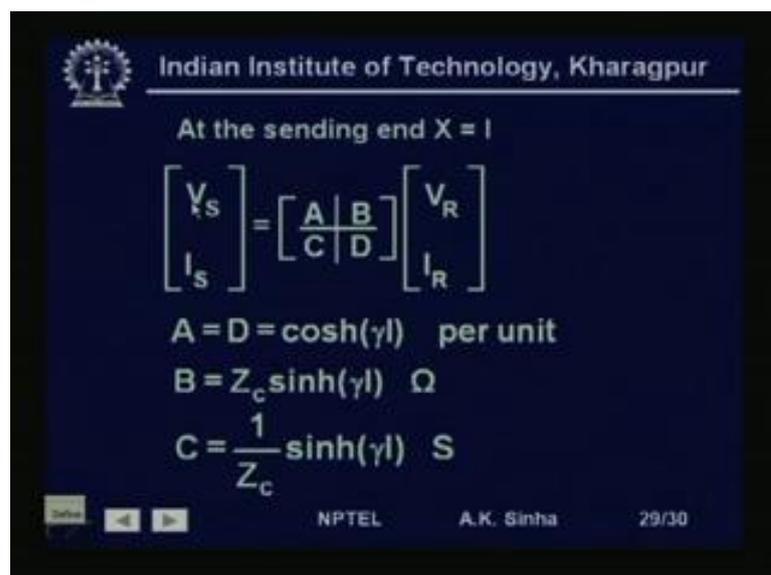
$$B(x) = Z_c \sinh(\gamma x) \ \Omega$$

$$C(x) = \frac{1}{Z_c} \sinh(\gamma x) \text{ S}$$

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So, this relationship can also be put in the form of sending end and receiving end voltage end current, where x is replaced by l .

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At the sending end $X = l$

$$\begin{bmatrix} V_S \\ I_S \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

$$A = D = \cosh(\gamma l) \text{ per unit}$$

$$B = Z_c \sinh(\gamma l) \ \Omega$$

$$C = \frac{1}{Z_c} \sinh(\gamma l) \text{ S}$$

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That is at sending end x is equal to l , therefore we can write this relationship as $V_S I_S$ is equal to $ABCD V_R I_R$. And where A is equal to D is \cos hyperbolic γl per unit, B is equal to $Z_c \sin$ hyperbolic γl , C is equal to 1 by $Z_c \sin$ hyperbolic γl in Siemens. So, again we have got the ABCD parameters for the transmission line and with for a long transmission line, the only thing here, what we are seeing is we have to use hyperbolic trigonometric functions.

That is ABCD parameters are in terms of hyperbolic trigonometric functions, not just Y and Z as we had for lumped parameter models. So, with this we stop today and we will be continuing this line long transmission line model. In the next class, where we will talk about the velocity of propagation the speed, the phase angles and another things. Also we will talk about regulation and efficiency of transmission lines.

Thank you.

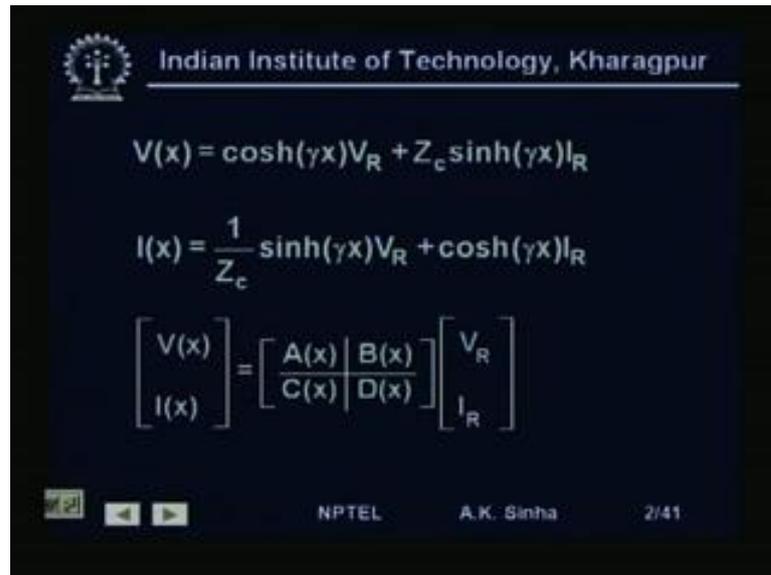
Preview of Next Lecture

Lecture - 08

Transmission Line Modeling Long Line (Contd.)

Welcome to lesson 8 on Power System Analysis. This lesson is a continuation on Transmission Line Modeling, specially Modeling of the Long Line. Now, if you remember in the previous lesson that is lesson 7, we talked about distributed parameter model for long transmission lines. That is lines which are longer than 250 kilometers.

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The slide displays the following equations and matrix representation for a transmission line:

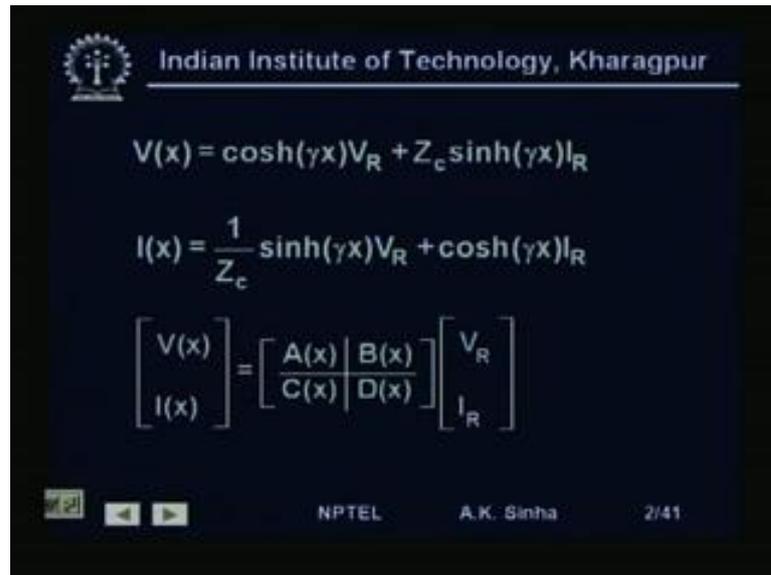
$$V(x) = \cosh(\gamma x)V_R + Z_c \sinh(\gamma x)I_R$$
$$I(x) = \frac{1}{Z_c} \sinh(\gamma x)V_R + \cosh(\gamma x)I_R$$
$$\begin{bmatrix} V(x) \\ I(x) \end{bmatrix} = \begin{bmatrix} A(x) & B(x) \\ C(x) & D(x) \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

At the bottom of the slide, there are navigation icons (back, forward, search), the text "NPTEL A.K. Sinha", and the slide number "2/41".

For these lines we said that the line voltage $V(x)$, at a distance x from the receiving end is given by $\cosh(\gamma x)V_R + Z_c \sinh(\gamma x)I_R$, where γ is called the propagation constant, x is the distance of the point from the receiving end, V_R and I_R are the voltages and currents at the receiving end. Similarly, the current at a distance x from the receiving end, $I(x)$ is equal to $\frac{1}{Z_c} \sinh(\gamma x)V_R + \cosh(\gamma x)I_R$, where Z_c is the characteristic impedance of the transmission line.

And it is given by square root of Z by Y , where Z is the series impedance of the transmission line per unit length. And Y is the shunt admittance of the transmission line per unit length. As we see these models these equations involve hyperbolic functions. Now, since we have been writing all these transmission line equations, in terms of ABCD parameters, we can write this equation also in those terms. So, in matrix form we can use this as $V(x), I(x)$ is equal to $A(x) B(x) C(x) D(x)$ and V_R and I_R this form, where this relationship is in terms of voltage and current at any distance x from the receiving end.

(Refer Slide Time: 54:43)



The slide displays the following equations and matrix representation for a transmission line:

$$V(x) = \cosh(\gamma x)V_R + Z_c \sinh(\gamma x)I_R$$
$$I(x) = \frac{1}{Z_c} \sinh(\gamma x)V_R + \cosh(\gamma x)I_R$$
$$\begin{bmatrix} V(x) \\ I(x) \end{bmatrix} = \begin{bmatrix} A(x) & B(x) \\ C(x) & D(x) \end{bmatrix} \begin{bmatrix} V_R \\ I_R \end{bmatrix}$$

At the bottom of the slide, there is a navigation bar with the text "NPTEL A.K. Sinha 2/41".

Now, here if you see $A(x)$ is equal to $D(x)$ and that is equal to \cos hyperbolic γx in per unit. That is if you see this relationship this is your $A(x)$ and this is your $B(x)$, this is $C(x)$ and this is $D(x)$. So, $A(x)$ is equal to $D(x)$ is equal to \cos hyperbolic γx , $B(x)$ is equal to $Z_c \sin$ hyperbolic γx , $V(x)$ is equal to $Z_c \sin$ hyperbolic γx . $C(x)$ is equal to $1/Z_c \sin$ hyperbolic γx and $D(x)$ as we have all ready seen is equal to $A(x)$.

So, this is the model, where A and D are basically dimensionless, B has a dimension of impedance and C has a dimension of Siemens that is admittance. Now, normally we are interested only in the terminal conditions. That is the sending end voltages end currents and the receiving end voltages end current. Rather than, the voltage end current at any intermediate point on the transmission line. Therefore, we can find out the voltage end currents at the sending end, in terms of voltage end current at the receiving end by substituting x is equal to l , where l is total line length.

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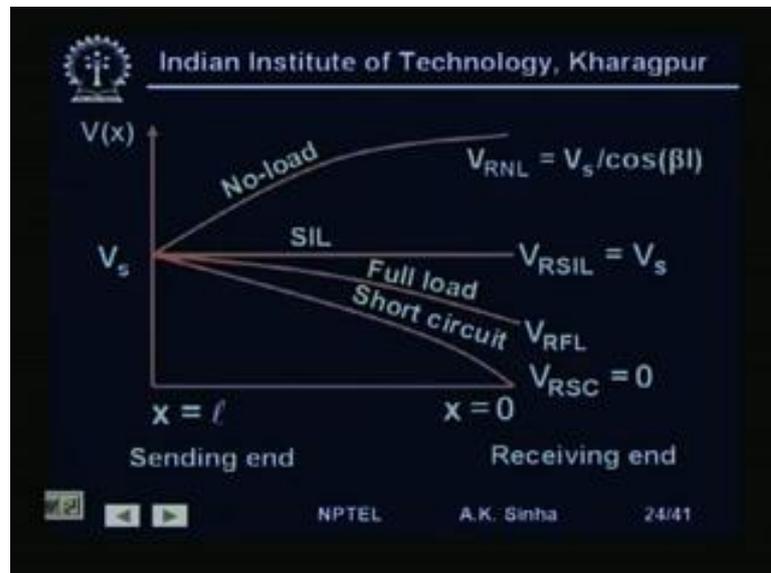
$$S(x) = P(x) + jQ(x) = V(x)I^*(x)$$
$$= \left(e^{j\beta x} V_R \right) \left(\frac{e^{j\beta x} V_R}{Z_c} \right)^*$$
$$= \frac{|V_R|^2}{Z_c}$$

Real Power flow along the line is constant and reactive power flow is zero

NPTEL A.K. Sinha 22/41

See this relationship, the voltage that is the power flowing is proportional to square of the voltage. So, if you are doubling the voltage, you are able to transmit 4 times the power and so on. Now, I will show you the characteristics of this long lossless transmission line.

(Refer Slide Time: 56:47)



When we have a surge impedance loading. The voltage across the line all along from the sending end to receiving end is going to be same. If the line is unloaded, then the sending end voltage.

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V_{rated} (KV)	$Z_c = \sqrt{L/C}$ Ω	$SIL = V_{rated}^2 / Z_c$ (MW)
230	380	140
345	285	420
500	250	1000
765	257	2280

$$V_{NL}(x) = (\cos \beta x) V_{RNL}$$
$$V_{SC}(x) = (Z_c \sin \beta x) I_{RSC}$$

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V S or V no load at any distance x which we can write as distance l is equal to cos beta x or cos beta l, for sending end voltage into V RNL. Now, in this case we find that the sending end voltage is going to be less than the receiving end voltage. And if we keep the sending end voltage as 1 per unit, then what we find at receiving end voltage is going to be much higher.

Thank you and in the next class, we will take up some problems on transmission lines.